



KØBENHAVNS
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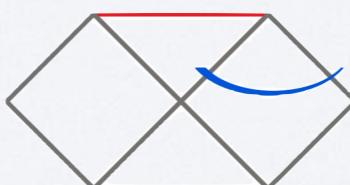


VILLUM FONDEN




HYPERBOLOIDAL APPROACH ON BH SPACETIMES: THE MINIMAL GAUGE

ArXiv: 2307.15735 - Royal Society Phil.Trans.A
“At the interface of asymptotic conformal methods and analysis in GR”



Rodrigo Panosso Macedo

OUTLINE

- Motivation
- The development of the minimal gauge
- Practical “recipe”: change signs in the tortoise coordinate
- Conclusion
- Projects

MOTIVATION

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(Umberto Eco - Foucault's Pendulum)

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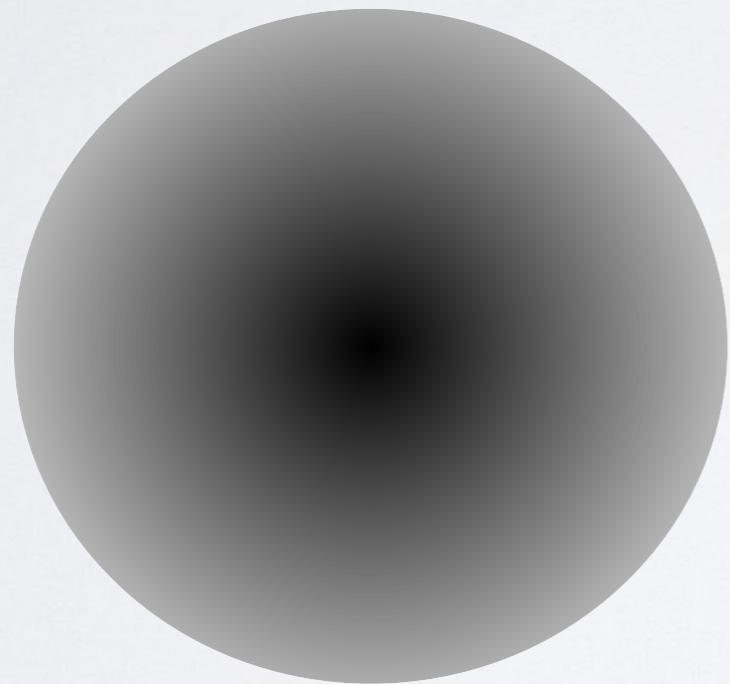
“That's why the Pendulum disturbs me. It promises the infinity, but where to put the infinity is left to me. So it isn't enough to worship the Pendulum; you still have to make a decision, you have to find the best point for it. And yet...”

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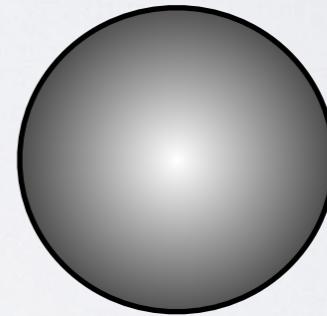
“That’s why the **Conformal Factor** disturbs me. It promises the infinity, but where to put the infinity is left to us. So it isn’t enough to worship the **Conformal Factor**; we still have to make a decision, we have to find the best point for it. And yet...”

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Physical manifold (\mathcal{M}, g)

$$g = \Omega^{-2} \bar{g}$$
A thick black arrow pointing from the large circle to the small circle, indicating a transformation or projection.



Unphysical manifold $(\bar{\mathcal{M}}, \bar{g})$

Boundary: $\partial \bar{\mathcal{M}}$ ($\Omega = 0$)

Decision: coordinates for hyperboloidal slices in black-hole spacetimes

BLACK HOLE GAUGES

- Schwarzschild (Boyer-Lindquist) coordinates: (t, r, θ, φ)
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Perturbation Theory

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Constraint Equations + Non-Linear Evolutions

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tortoise coordinate

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$$ds^2 = \Omega^{-2} d\bar{s}^2$$

$$\Omega = \frac{\sigma}{\lambda}$$

$$d\bar{s}^2 = \Xi(\sigma) [-p(\sigma)d\tau^2 + 2\gamma(\sigma)d\tau d\sigma + w(\sigma)d\sigma^2] + \rho(\sigma)^2 d\omega^2$$

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tortoi

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$$\Omega = \frac{\sigma}{\lambda}$$

Decision #0

$$d\bar{s}^2 = \Xi^2 d\sigma^2 + \Xi^2 r_*^2 d\omega^2$$

Use conformal factor as radial coordinate: $\sigma = \lambda \Omega$

HYPERBOLOIDAL METRIC

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Radial degrees of freedom

$\rho(\sigma)$: Conformal areal radius $\beta(\sigma) := \rho(\sigma) - \sigma\rho'(\sigma)$

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$$\mathcal{F}(\sigma) = \sqrt{a(r)b(r)}$$

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Horizons: $\mathcal{F}(\sigma_h) = 0$

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Fancy Metric: $\zeta(\sigma) \neq 1$

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$\bar{\eta}_{ab}$

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Time degrees of freedom

$$\gamma(\sigma) := -\lambda \frac{dH}{dr_*} = H'(\sigma)p(\sigma)$$

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Time degrees of freedom

$$\gamma(\sigma) := -\lambda \frac{dH}{dr_*} = H'(\sigma)p(\sigma) \quad w(\sigma) := \frac{1 - \gamma(\sigma)^2}{p(\sigma)}$$

HYPERBOLOIDAL GEOMETRY

- Compactified hyperboloidal coordinates: $(\tau, \sigma, \theta, \varphi)$

$$d\bar{s}^2 = \Xi(\sigma) [-p(\sigma)d\tau^2 + 2\gamma(\sigma)d\tau d\sigma + w(\sigma)d\sigma^2] + \rho(\sigma)^2 d\omega^2$$

Causal structure

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$$\bar{l}^a = \nu \left(\delta_\tau^a - \frac{1+\gamma}{w} \delta_\sigma^a \right)$$

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(Normal vector to
time surfaces)

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(Normal vector to
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Hyperboloidal Surface

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Hyperboloidal Surface

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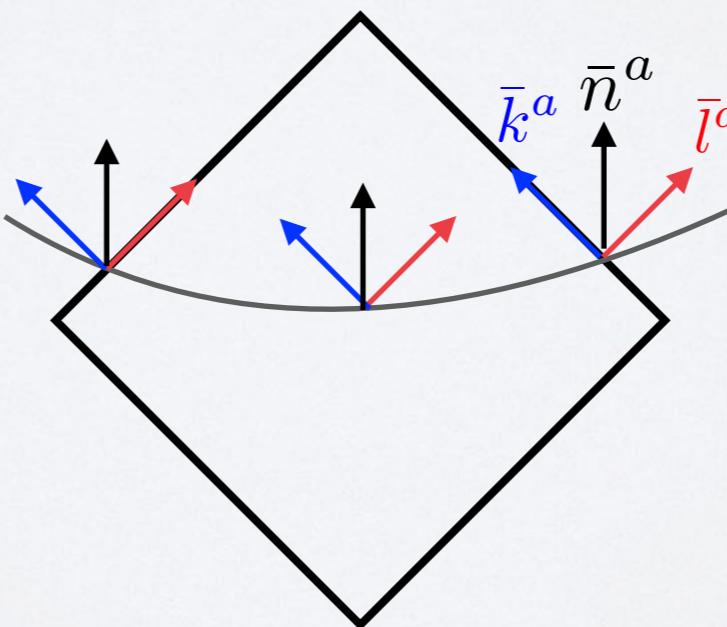
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Hyperboloidal Surface



SCHWARZSCHILD

$$ds^2 = -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2d\omega^2, \quad f(r) = 1 - \frac{r_h}{r}$$

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M. Ansorg

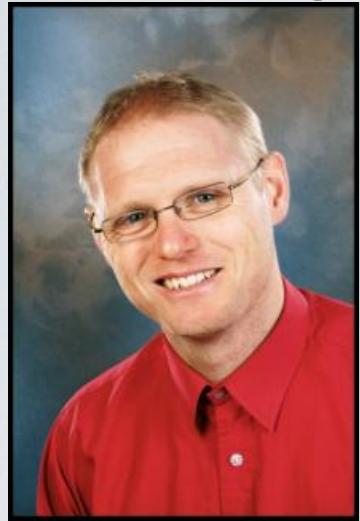


(★1970; †2016)

SCHWARZSCHILD

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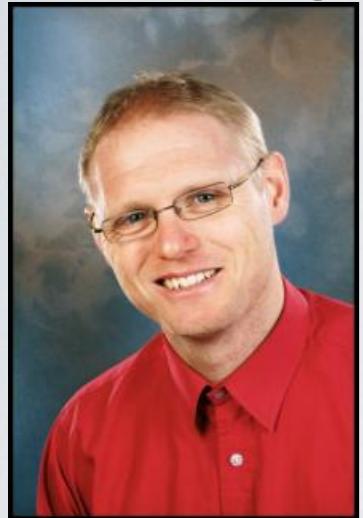
I. Ingoing null coordinates: $v = t - r_*$

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SCHWARZSCHILD

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SCHWARZSCHILD

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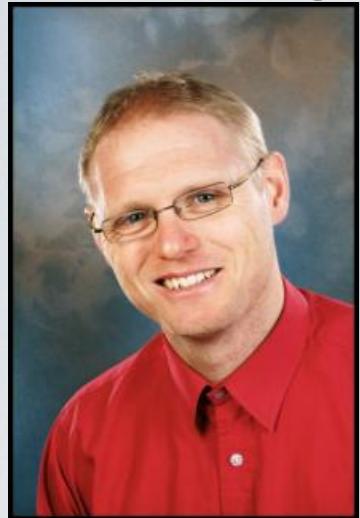
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SCHWARZSCHILD

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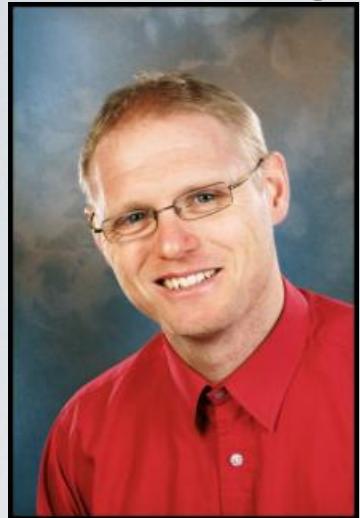
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2. Consider asymptotic behaviour of outgoing light rays

SCHWARZSCHILD

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M. Ansorg



(★1970; †2016)

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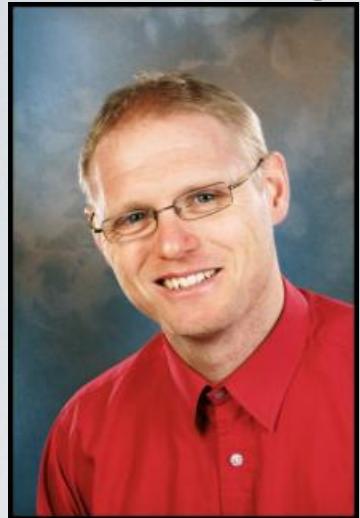
2. Consider asymptotic behaviour of outgoing light rays

$$\frac{dv}{dr} = \frac{2}{1 - r_h/r} \approx 2 \left(1 + \frac{r_h}{r} \right)$$

SCHWARZSCHILD

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M. Ansorg



(★1970; †2016)

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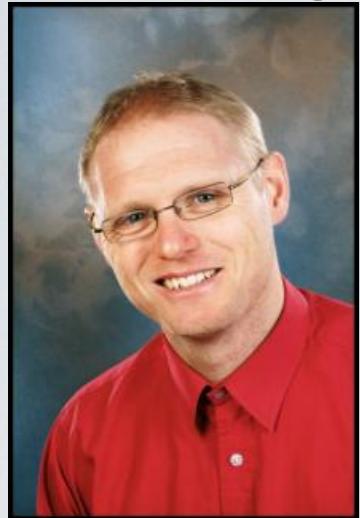
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$$v = 2r + 2r_h \ln \left(\frac{r}{r_h} \right) + C$$

SCHWARZSCHILD

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(★1970; †2016)

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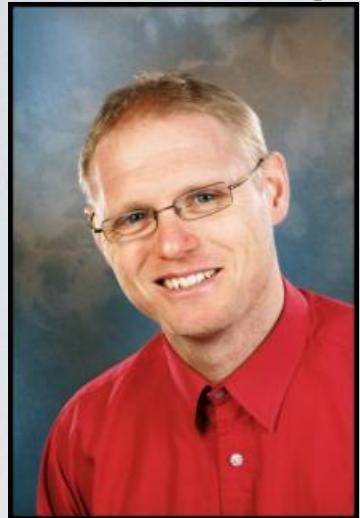
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3. Choose integration constant as time coordinate $C = \lambda\tau$

SCHWARZSCHILD

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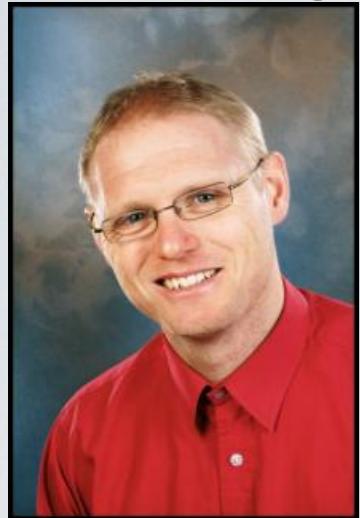
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SCHWARZSCHILD

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$$r = \frac{r_h}{\sigma}$$

SCHWARZSCHILD

$$d\bar{s}^2 = -\sigma^2(1-\sigma)d\tau^2 + \frac{2r_h}{\lambda}(1-2\sigma^2)d\tau d\sigma + \frac{4r_h^2}{\lambda^2}(1+\sigma)d\sigma + \frac{r_h^2}{\lambda^2}d\omega^2$$

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SCHWARZSCHILD

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SCHWARZSCHILD

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SCHWARZSCHILD

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SCHWARZSCHILD

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SCHWARZSCHILD

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Time degree of freedom $t = \lambda \left(\tau - H(\sigma) \right)$

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SCHWARZSCHILD

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$$\rho(\sigma) = \frac{r_h}{\lambda} \xrightarrow{\beta(\sigma) := \rho(\sigma) - \sigma\rho'(\sigma)} \beta(\sigma) = \frac{r_h}{\lambda}$$

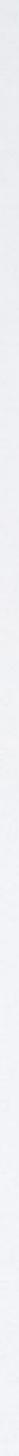
$$x(\sigma) = \frac{r_*(\sigma)}{\lambda} = \frac{r_h}{\lambda} \left(\frac{1}{\sigma} - \ln \sigma + \ln(1-\sigma) \right)$$

Time degree of freedom $t = \lambda(\tau - H(\sigma))$

$$H(\sigma) = \frac{r_h}{\lambda} \left(\frac{1}{\sigma} + \ln \sigma + \ln(1-\sigma) \right)$$

QUASINORMAL MODES

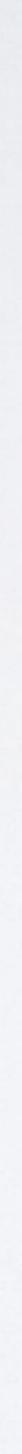
(Leaver 1985) $\lambda = r_h \Leftrightarrow 2M = 1$



QUASINORMAL MODES

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(t, r, θ, φ)

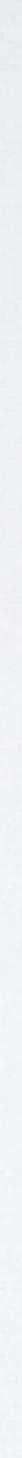


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$$(t, r, \theta, \varphi)$$

$$-\psi_{,tt} + \psi_{,r_*r_*} - V(r)\psi = 0$$



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$$\psi(t, r) = e^{-i\omega t}\phi(r)$$



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$$\phi(r) = \underbrace{r^{2i\omega}(r-1)^{-i\omega}}_{\text{BC}} \underbrace{e^{2i\omega(r-1)}}_{\text{regular}} \bar{\phi}_L(r)$$

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(Horizon)

$$\bar{\phi}_L(r) = \sum_{n=0}^{\infty} a_n \left(1 - \frac{1}{r}\right)^n$$

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Radial
transformation in
Ansorg's Recipe

Height Function in
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Take Away Message

(M. Ansorg, RPM 2016)

Leaver's regular field is the frequency-domain representation of the hyperboloidal perturbation field $\bar{\psi}(\tau, \sigma)$ in the minimal gauge

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Radial transformation in Ansorg's Recipe

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Decision #1

Follow Ansorg's Recipe

1. Ingoing null coordinates
2. Integrate asymptotic behaviour of outgoing light rays
3. Choose integration constant as time coordinate
4. Simple compactification

REISSNER-NORDSTRÖM

$$ds^2 = -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2d\omega^2, \quad f(r) = 1 - \frac{2M}{r} + \frac{Q^2}{r^2}$$

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“you still have to make a decision, you have to find the best point for it.

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I. Ingoing null coordinates: $v = t - r_*$

$$ds^2 = dv \left(-f(r)dv^2 + 2dr \right) + r^2d\omega^2$$

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1. Ingoing null coordinates: $v = t - r_*$

$$ds^2 = dv \left(-f(r)dv^2 + 2dr \right) + r^2d\omega^2$$

2. Consider asymptotic behaviour of outgoing light rays

$$\frac{dv}{dr} = \frac{2}{f(r)} \approx 2 \left(1 + \frac{2M}{r} \right)$$

$$v = 2r + 4M \ln \left(\frac{r}{r_h} \right) + C$$

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$$r = \frac{r_h}{\sigma}$$

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“you still have to make a decision, you have to find the best point for it. And yet...”

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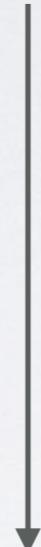
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Radial
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Height Function in
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Radial transformation in the Ansorg's Recipe Height Function in the Ansorg's Recipe

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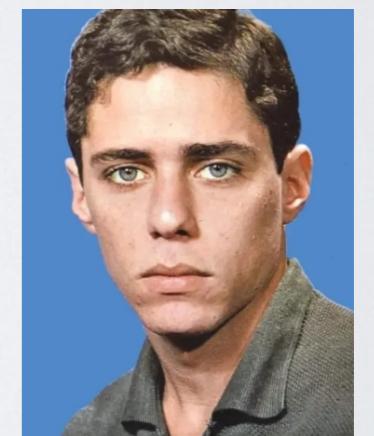
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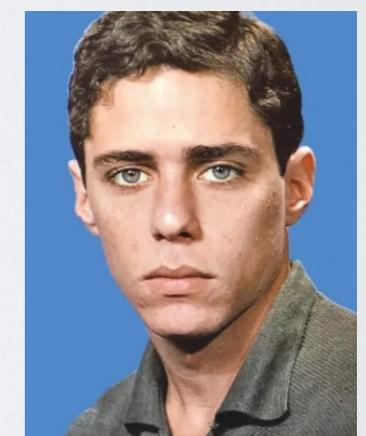
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THEY DON'T AGREE



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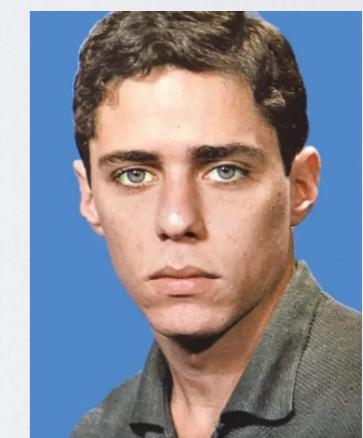
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LEAVER EXPRESSION'S SEEKS SIMPLER



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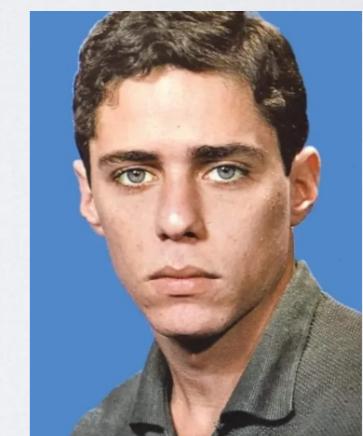
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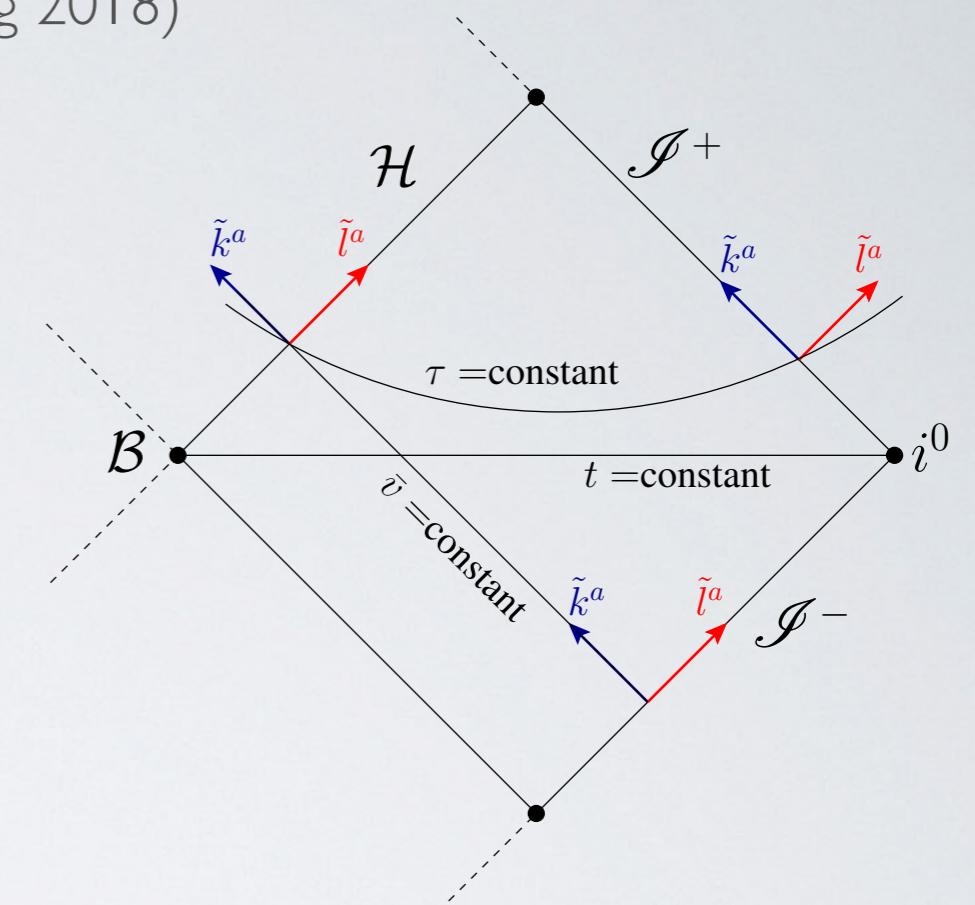
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MINIMAL GAUGE

(RPM, J.L. Jaramillo, M. Ansorg 2018)



4. Simple compactification

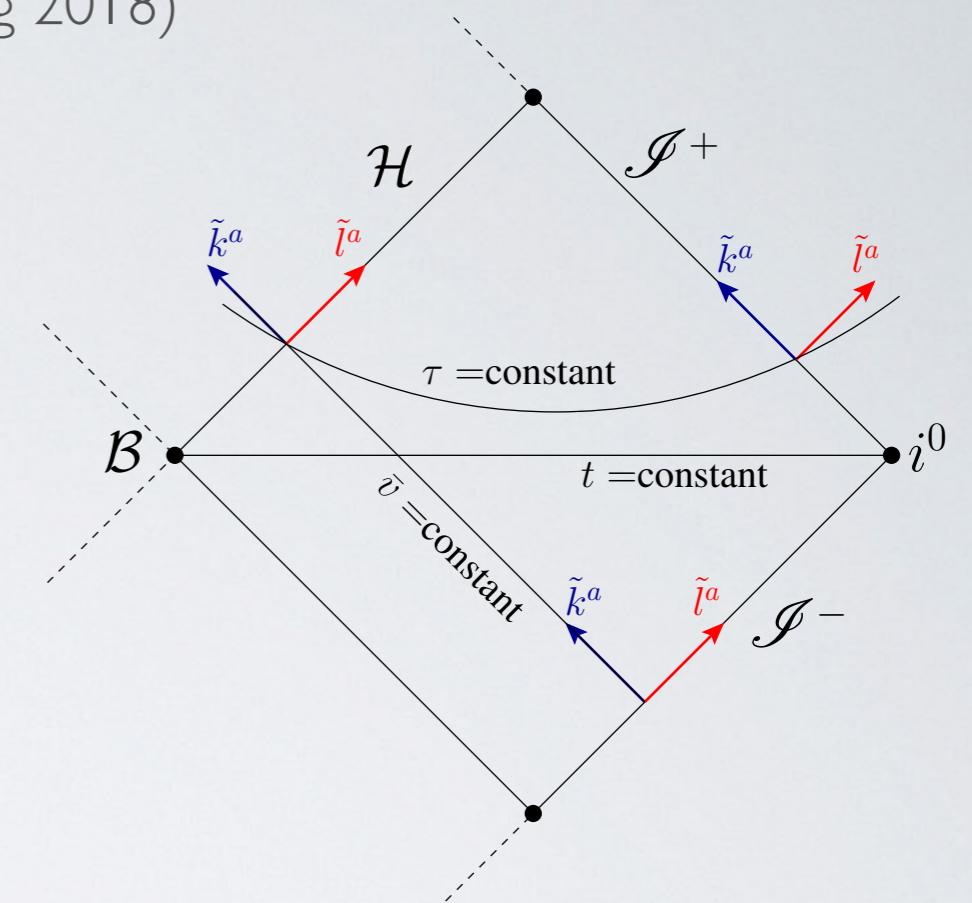


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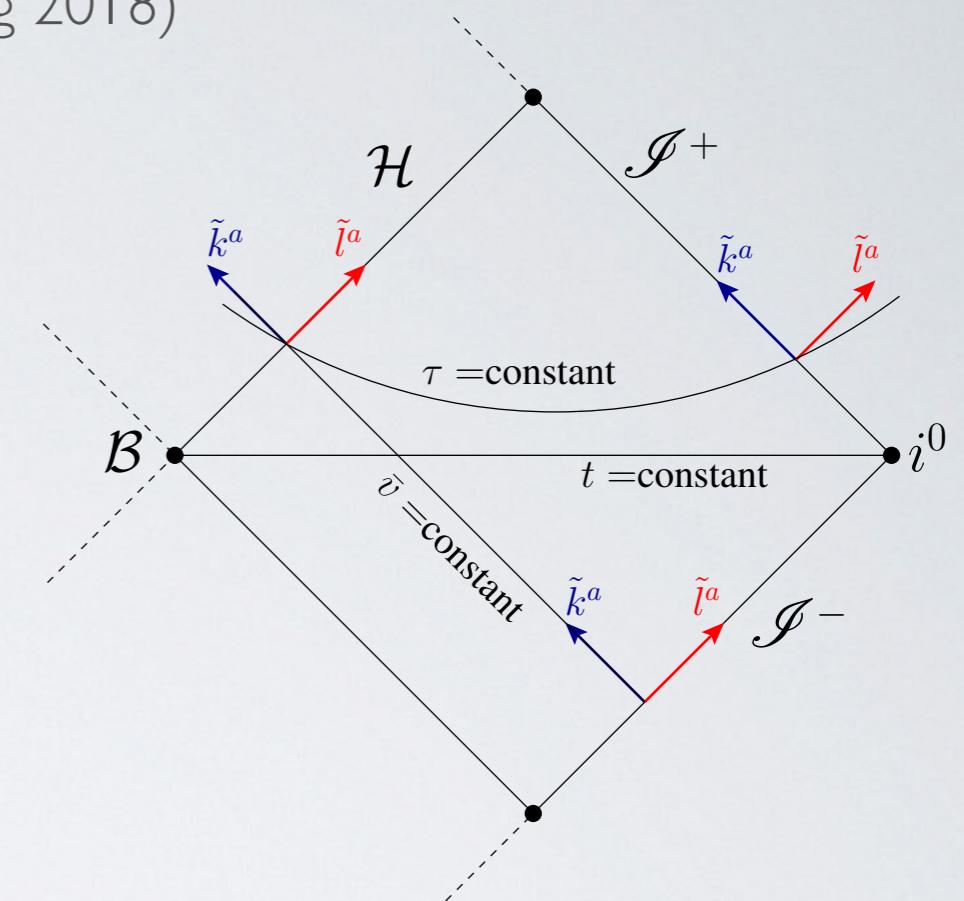
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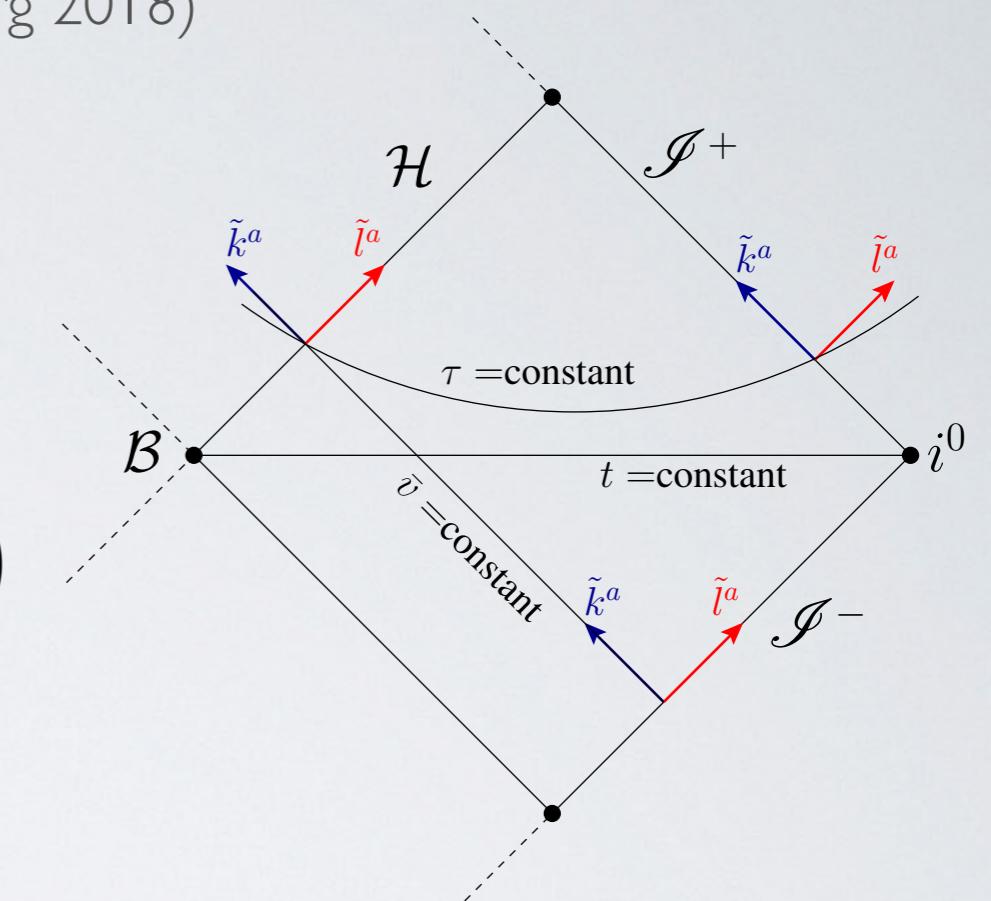
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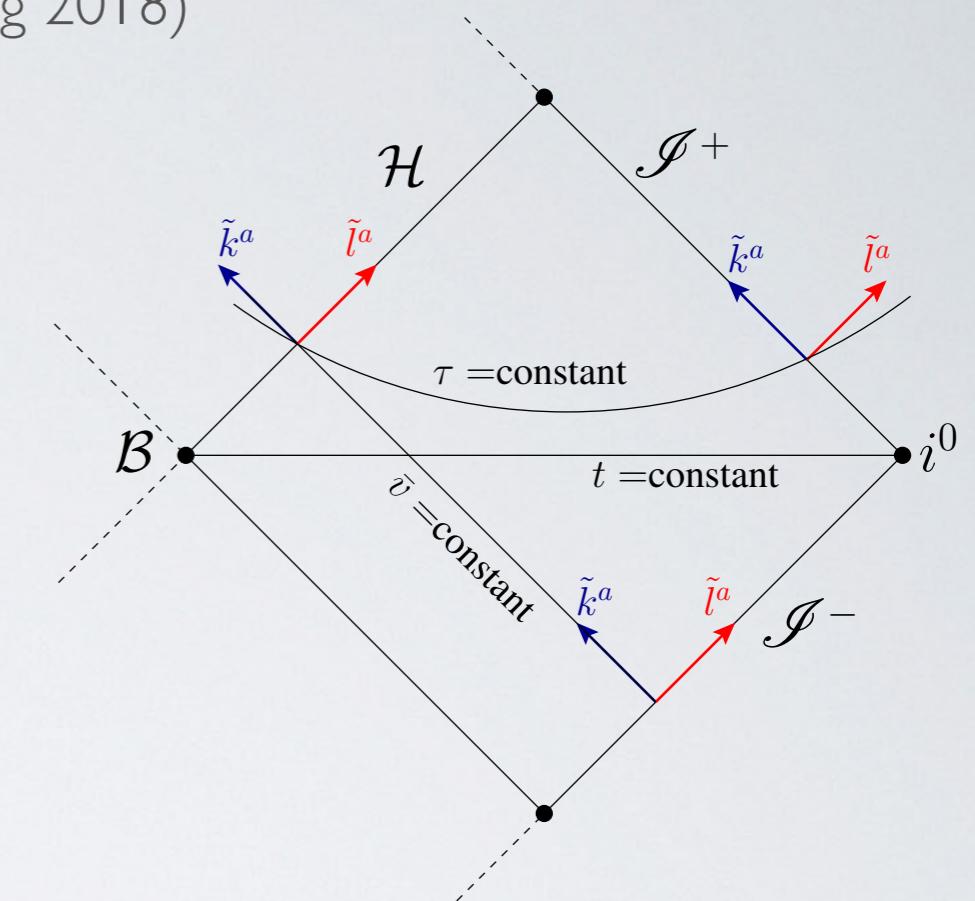
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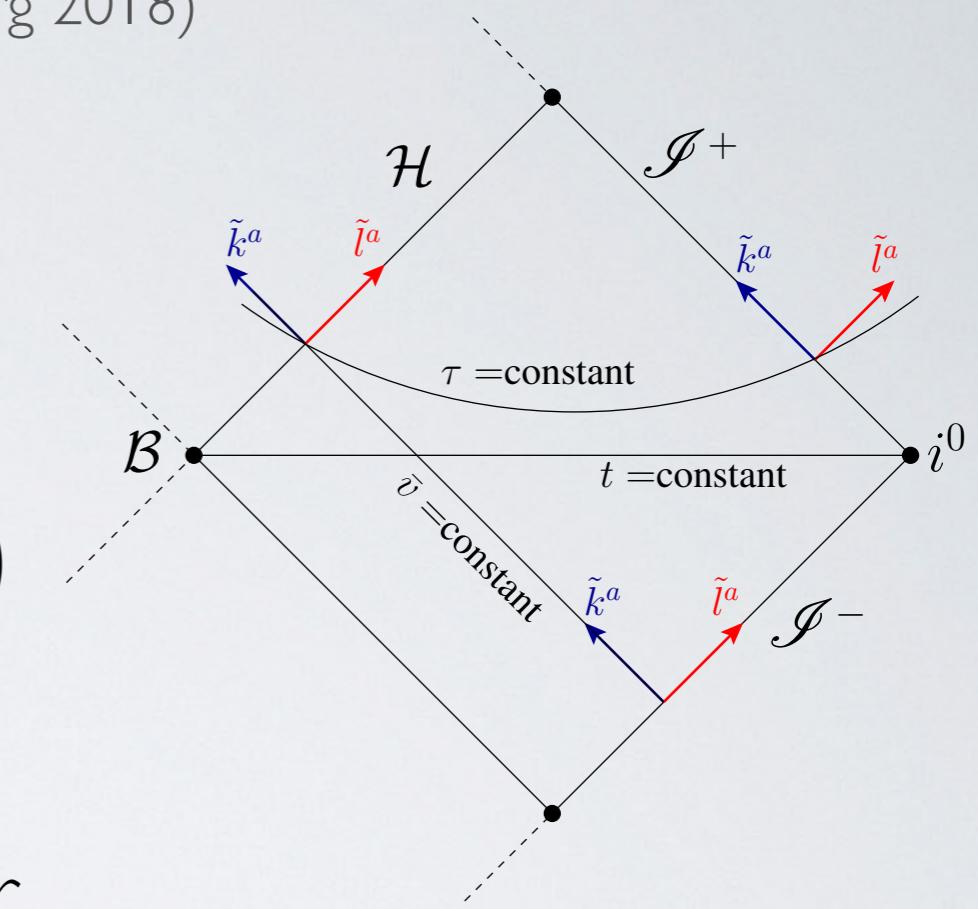
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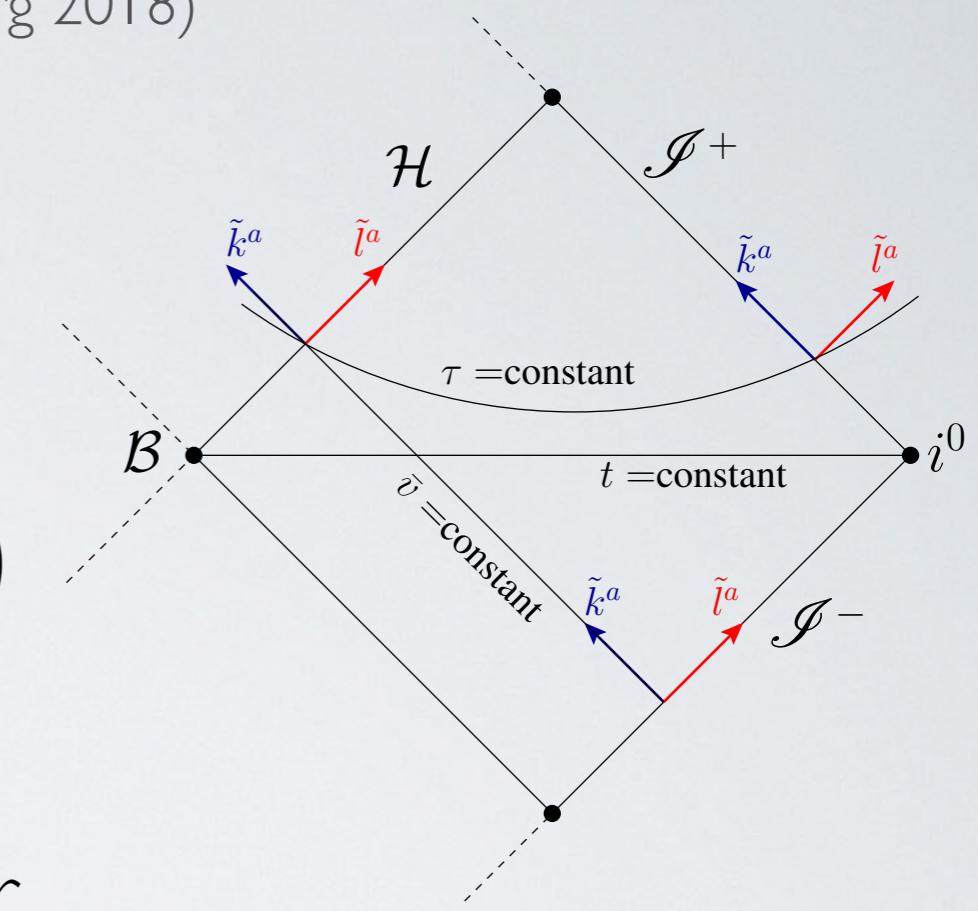
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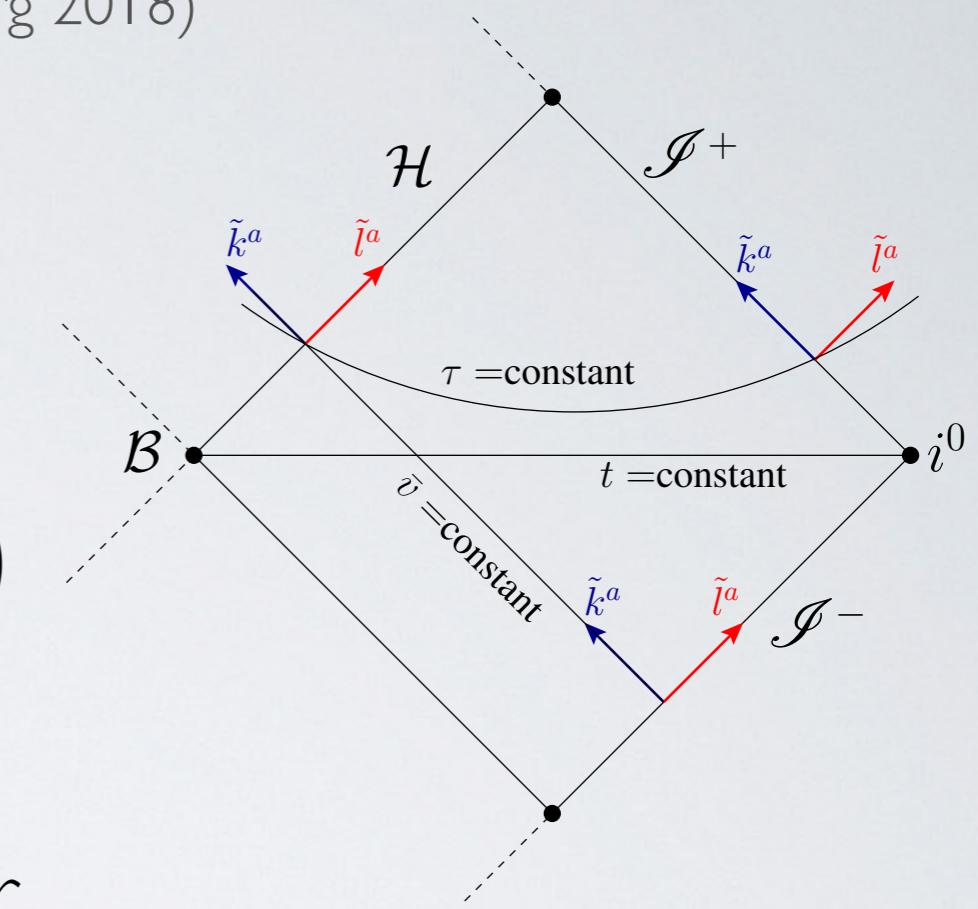
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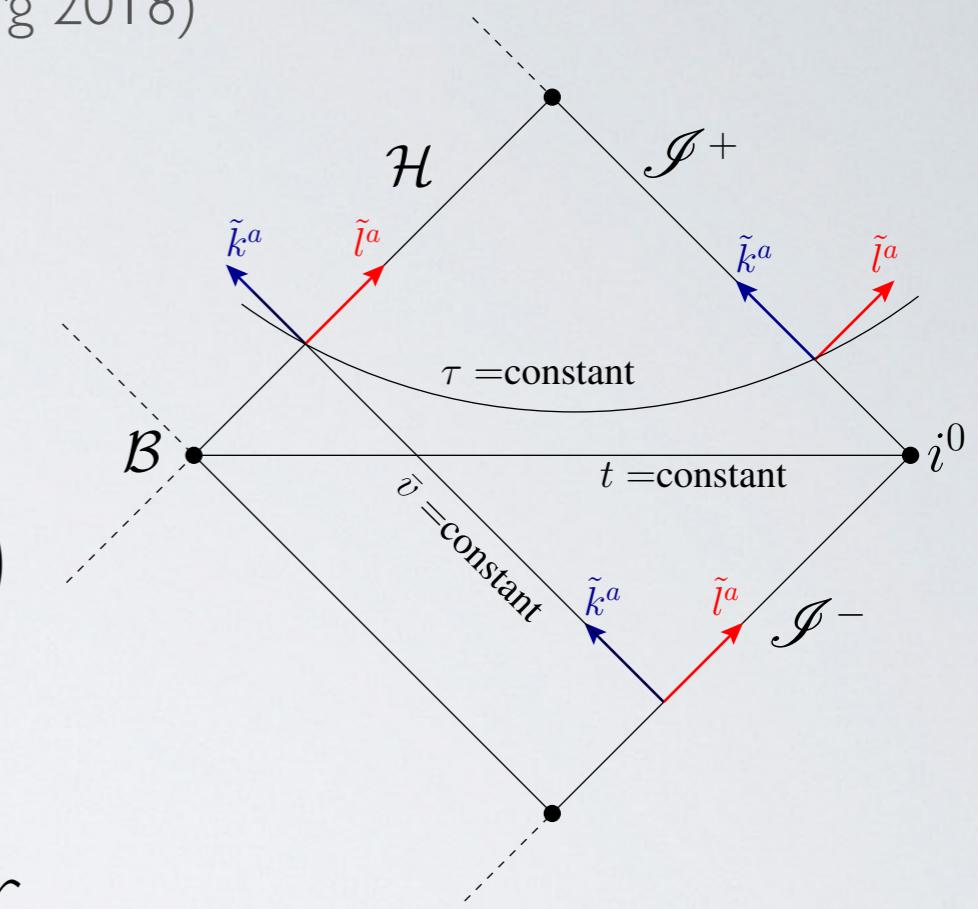
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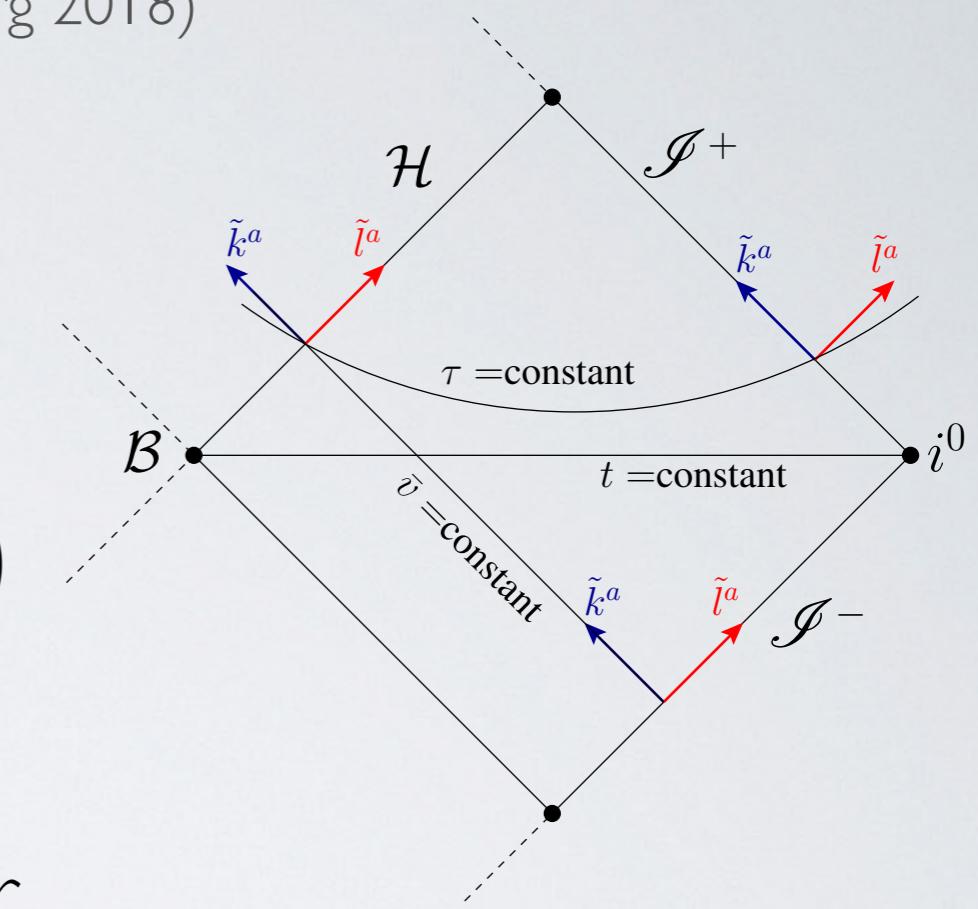
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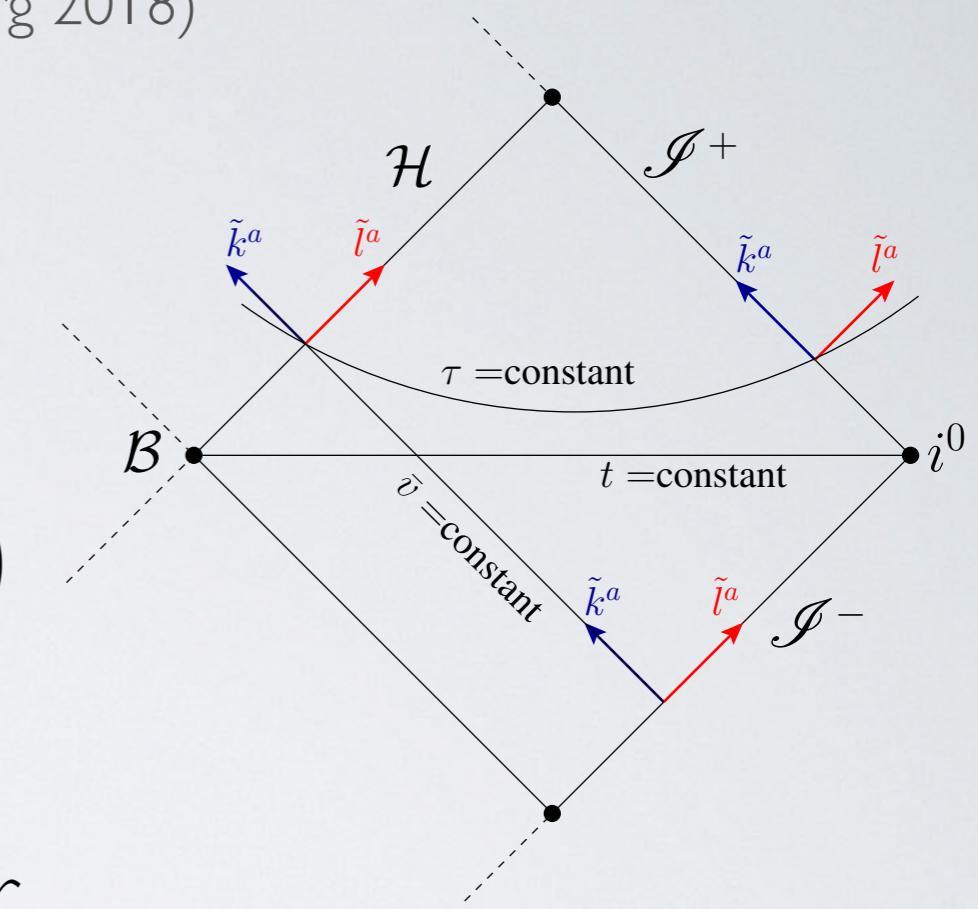
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REISSNER-NORDSTRÖM

(And Kerr - RPM 2019)

Minimal Gauge: reduce the gauge freedom to the most simple case

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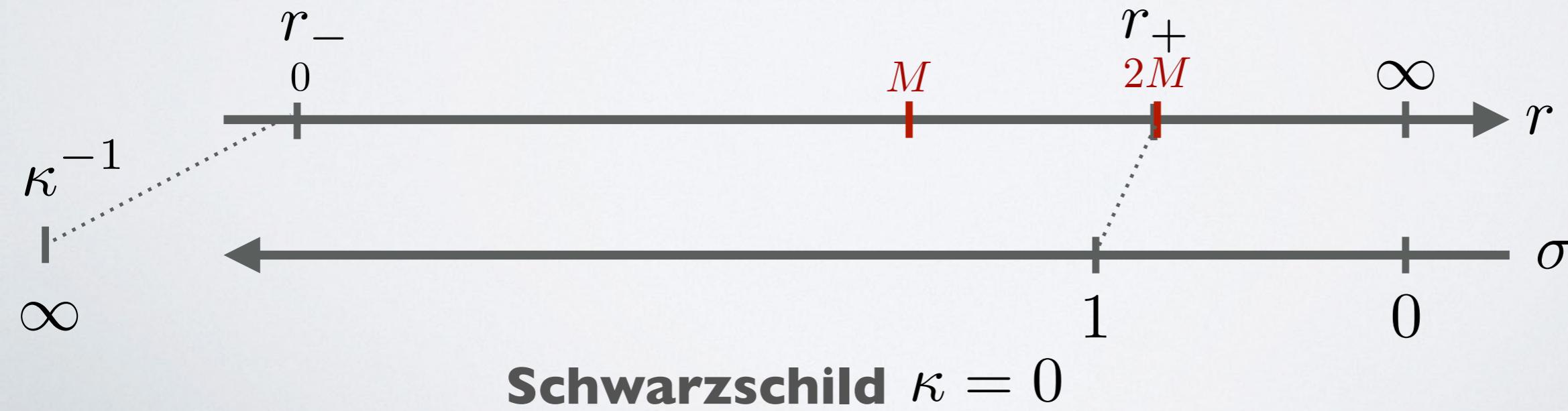
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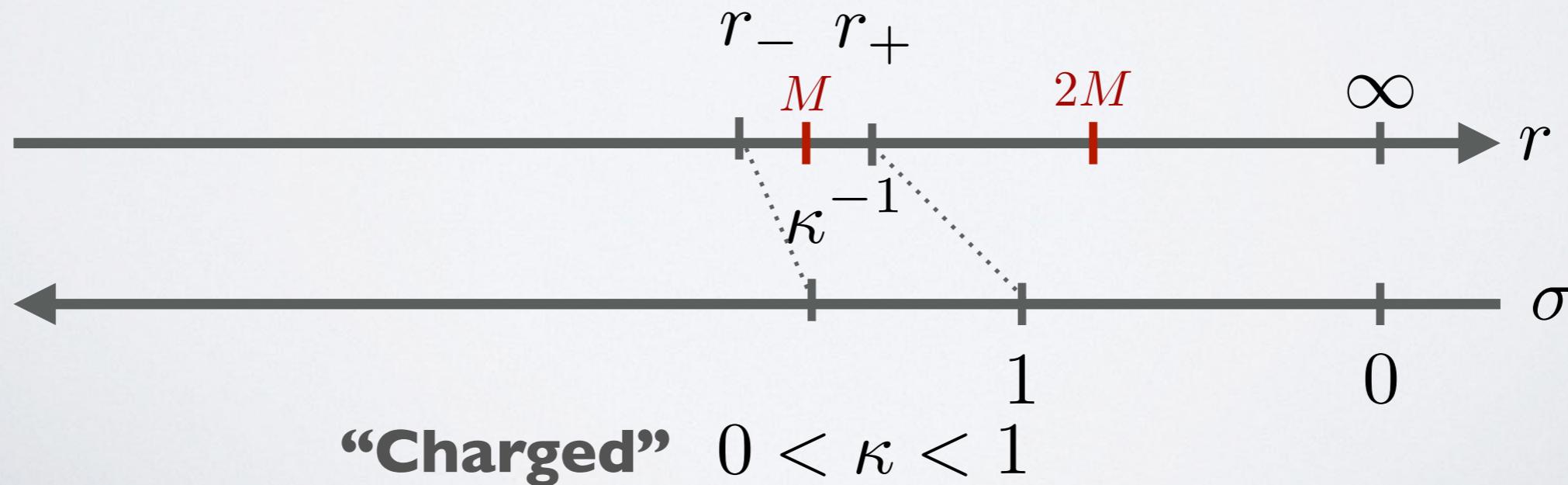
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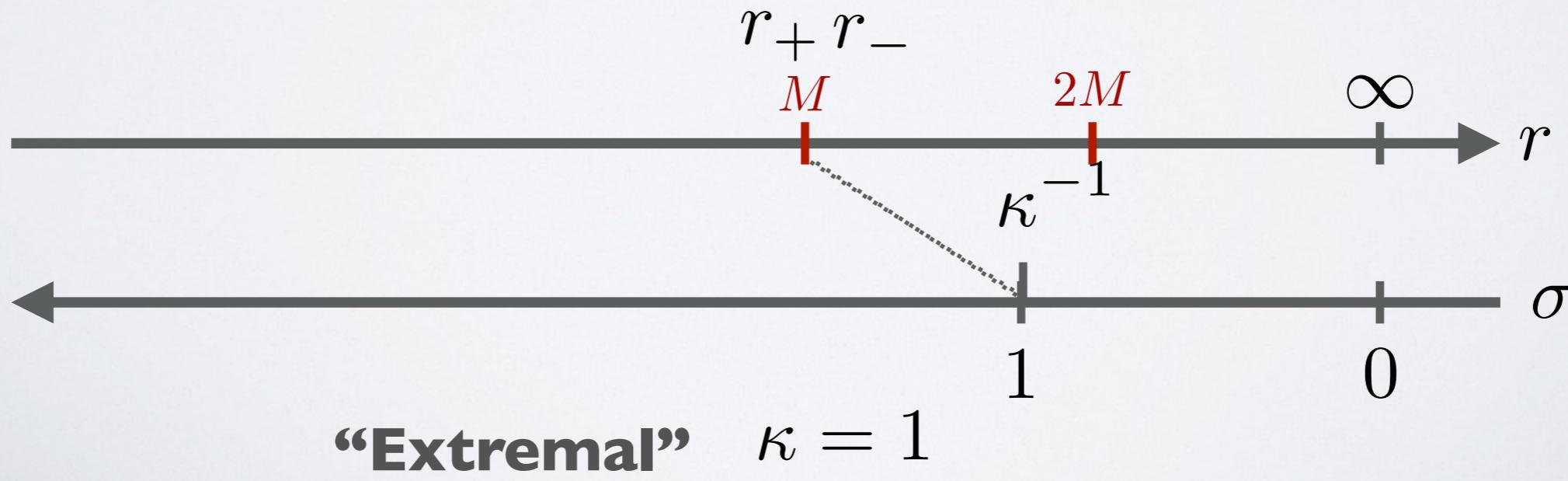
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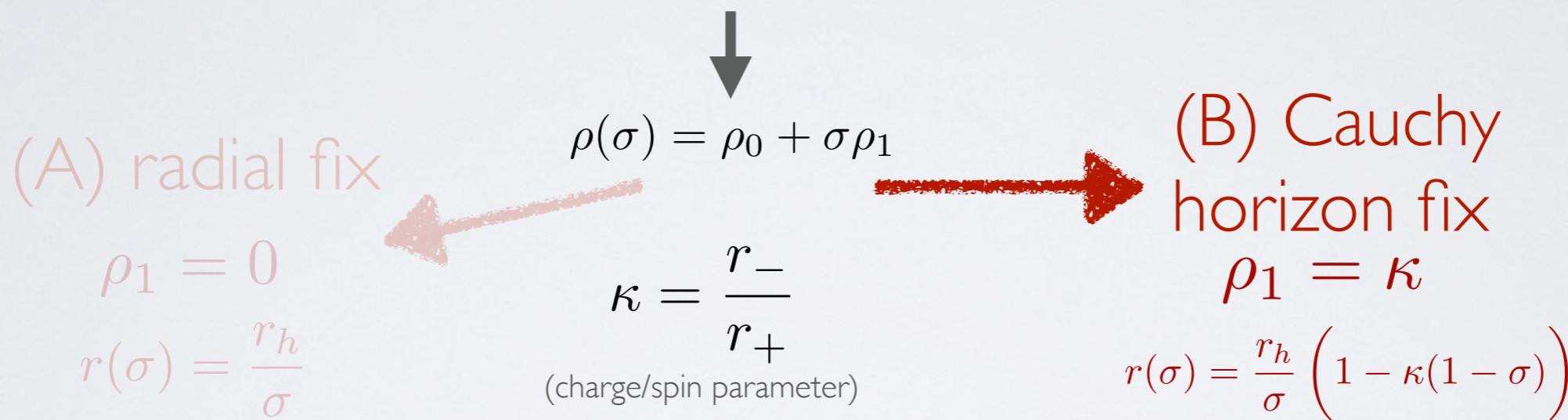
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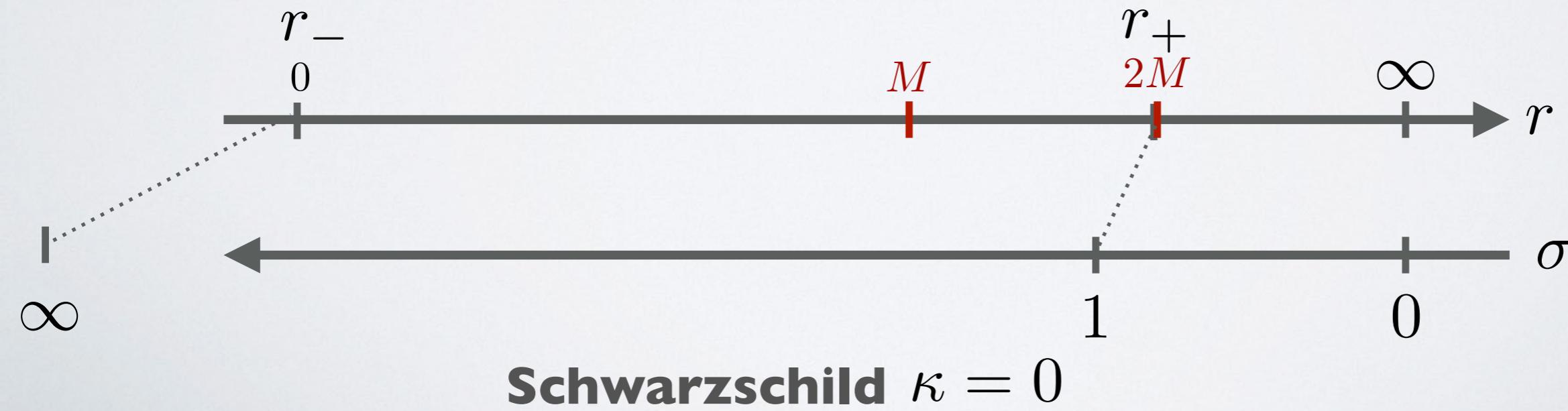
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(B) Cauchy horizon fix

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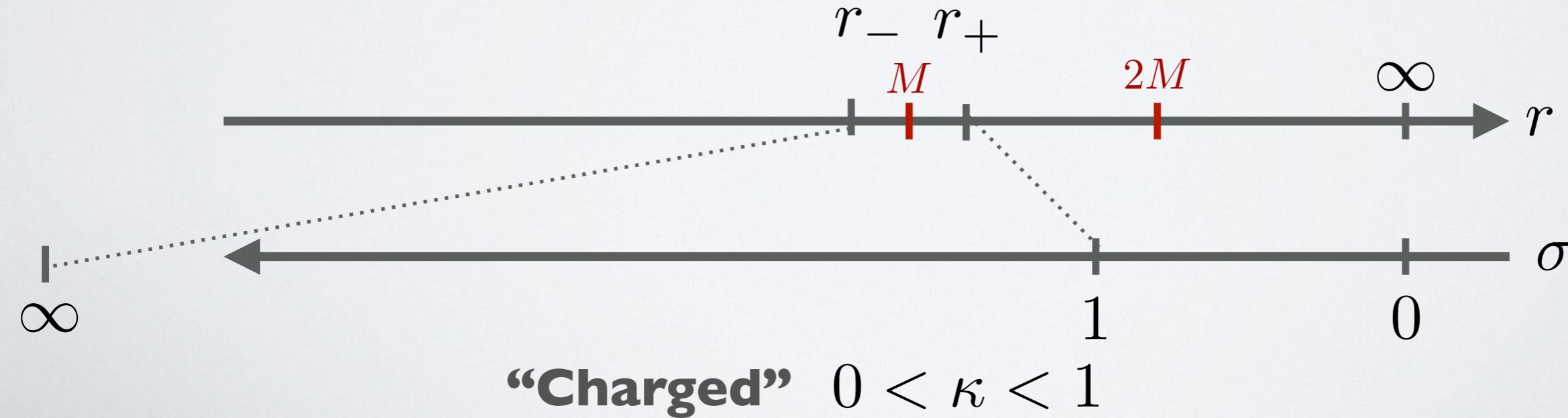
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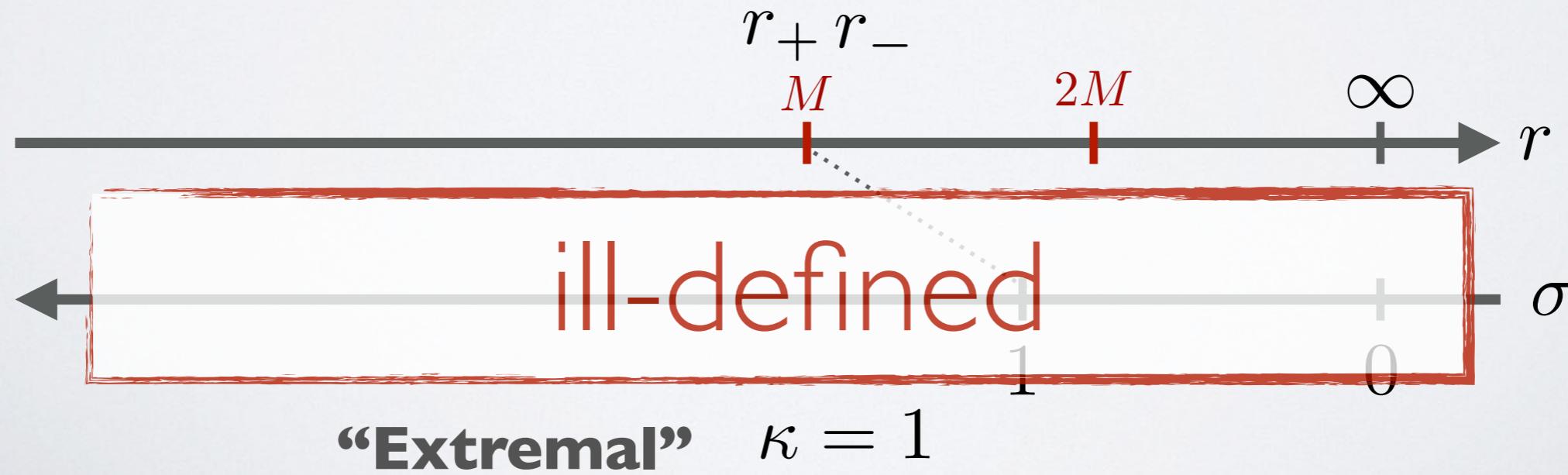
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REISSNER-NORDSTRÖM

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REISSNER-NORDSTRÖM

(B) Cauchy horizon fixing: the extremal limit

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The Robinson-Bertotti metric (near-horizon geometry): $AdS_2 \times S^2$

B. Bertotti, Phys. Rev 116 1331 (1959)

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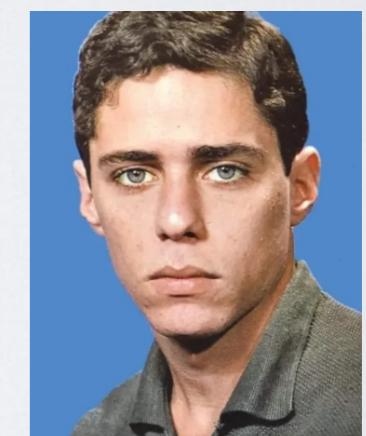
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Leaver's regular field is the frequency-domain
representation of hyperboloidal field in minimal gauge



Leaver's approach fails in a direct limit to the extremal case, because spacetime representation corresponds to discontinuous transition into the near-horizon geometry

BLACK HOLE + MATTER HALO

$$ds^2 = -a(r)dt^2 + \frac{dr^2}{b(r)} + r^2d\omega^2$$
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(V. Cardoso, K. Destounis, RPM,
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“you still have to make a decision, you have to find the best point for it. And yet...”

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$$d\bar{s}^2 = \Xi(\sigma) \left[-p(\sigma) d\tau^2 + 2\gamma(\sigma) d\tau d\sigma + w(\sigma) d\sigma^2 \right] + \frac{r_h^2}{\lambda^2} d\omega^2$$

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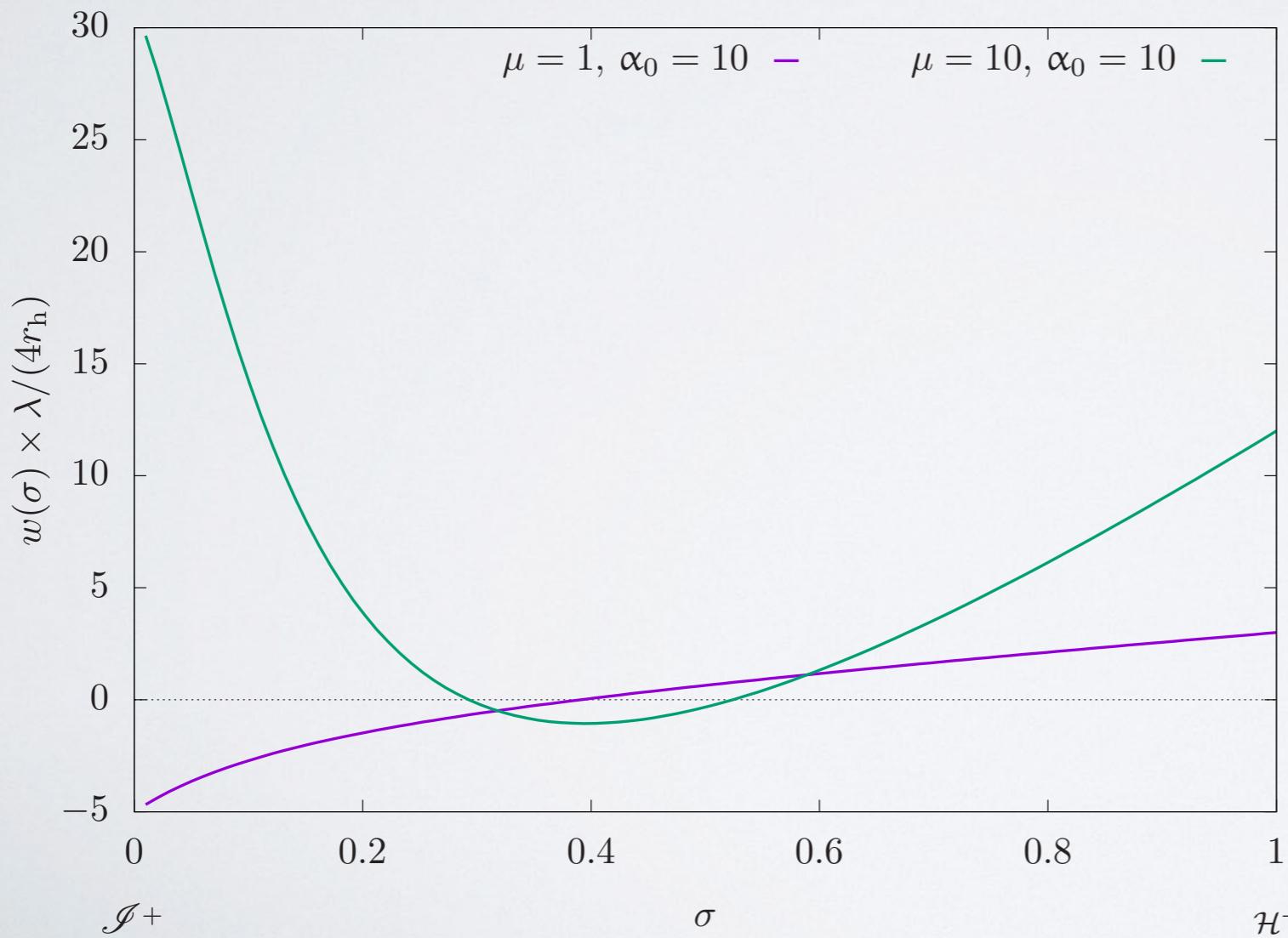
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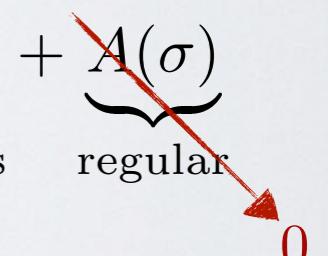
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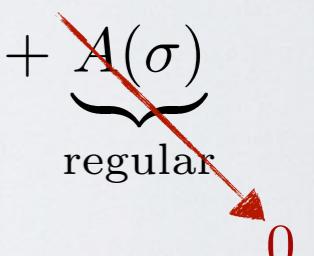
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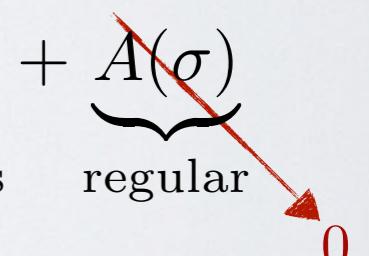
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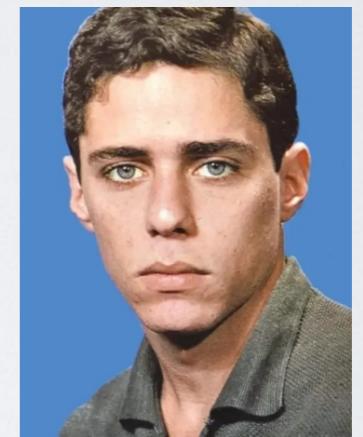
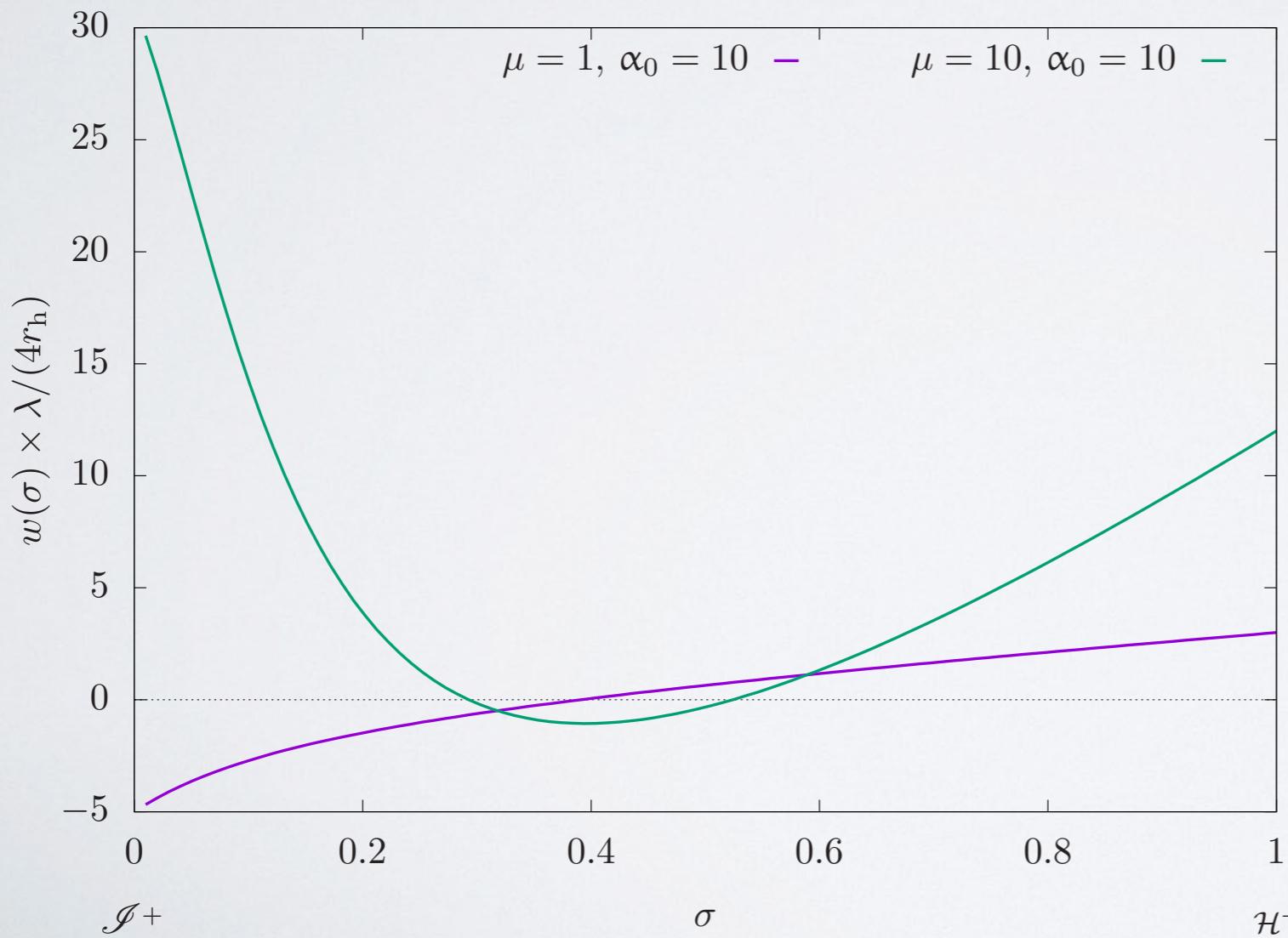
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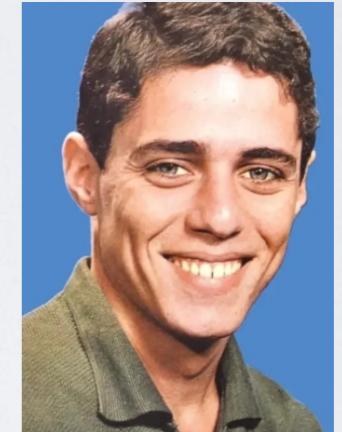
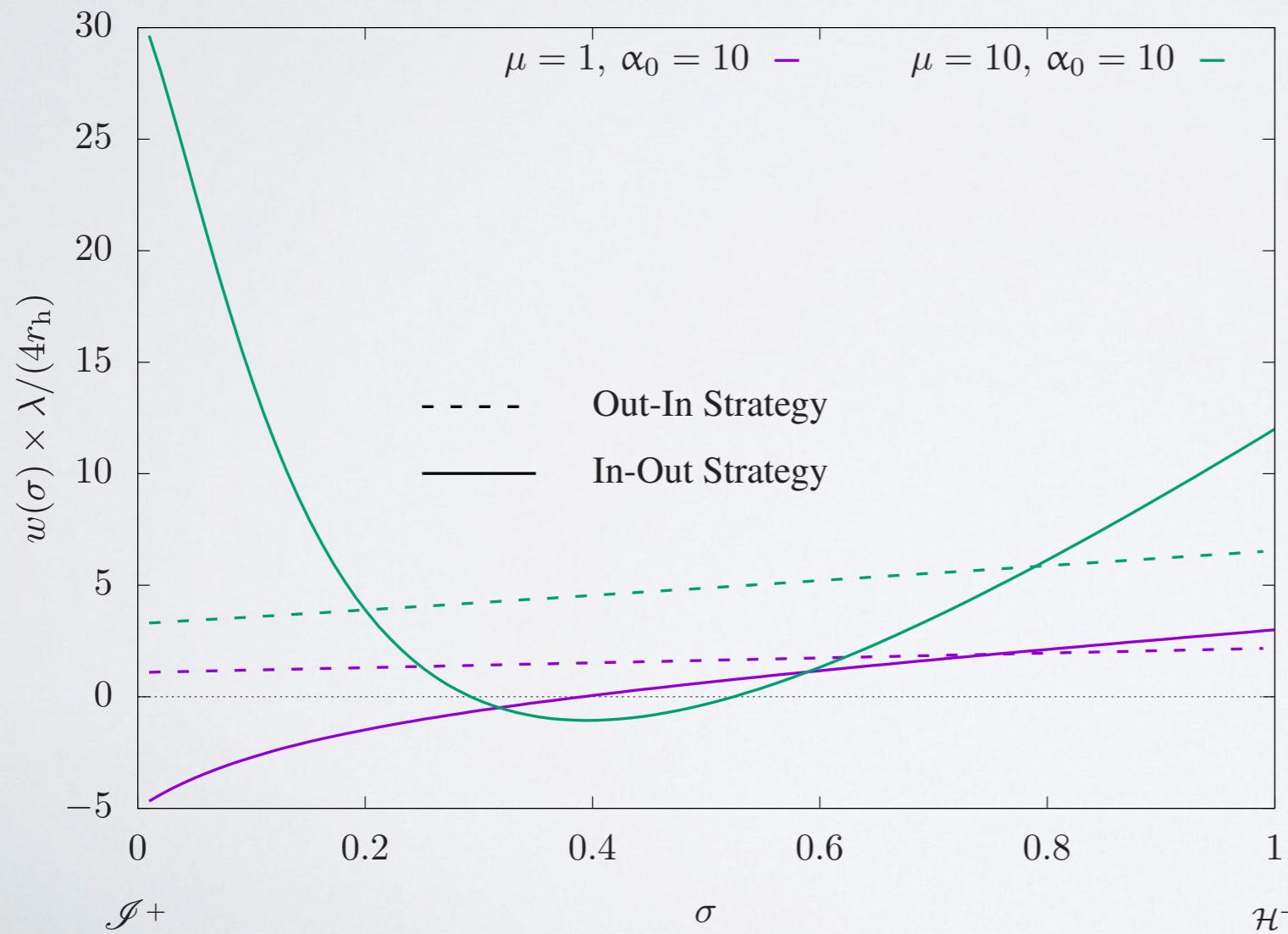


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(ArXiv 2307.15735)

MINIMAL GAUGE

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If $x_{\text{reg}}(\sigma) \neq \text{constant}$ then out-in strategy is best

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- Out-In Strategy: $H_{\text{out-in}}(\sigma) = \sum_{i=0}^{N_h} x_{h_i}(\sigma) - x_0(\sigma) - x_{\text{reg}}(\sigma)$

If $x_{\text{reg}}(\sigma) \neq \text{constant}$ then out-in strategy is best

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- Historical review on the conceptual framework behind the so-called hyperboloidal minimal gauge for black-hole spacetimes
- The hyperboloidal framework in the minimal gauge provides a geometrical (spacetime) interpretation for methods employed in the QNM literature
- Minimal gauge height function follows from just flipping signs in the tortoise coordinate
- There always exists a well-defined hyperboloidal foliation with the out-in strategy (empirical observation; no formal proof)
- Minimal gauge conceptual construction is based on coordinate change: there is no covariant classification to define it.

PROJECTS

Perturbation theory and Gravitational Wave astronomy

Second order: Self-force and ring-down
(e.g. regularity of source terms)

Environmental effects: coupling between waves with different characteristic speeds (regularity in time and frequency domain)

$$-\psi_{,tt} + \psi_{,r_* r_*} - V_\psi(r)\psi + C_\psi(r)\rho = 0$$

$$-v^{-2}\rho_{,tt} + \rho_{,r_* r_*} - V_\rho(r)\rho + C_\rho(r)\phi = 0$$

- Black-hole perturbation theory in the (regular) conformal set up:
Derive Regge-Wheller-Zerrili and/or Bardeen-Press-Teukolsky equations *directly* from the conformal equations?
(So far, using coordinate coordinate transformation, conformal re-scalings...)

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- Linear Regime $\square_o \chi^{(1)} = \mathcal{S}^{(1)}$

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PROJECTS

Perturbation theory and Gravitational Wave astronomy

Second order: Self-force and ring-down
(e.g. regularity of source terms)

- Linear Regime
- High order terms

$$\square_o \chi^{(1)} = \mathcal{S}^{(1)}$$

$$\square_o \chi^{(2)} = \mathcal{S}^{(2)}(\chi^{(1)}, \partial \chi^{(1)})$$

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